



















Attenuation $n = \mu N\Delta x \text{ photons lost per unit length}$ $\mu = \frac{n/N}{\Delta x} \text{ fraction of photons lost per unit length}$ $\Delta N = -n \longrightarrow \frac{dN}{dx} = -\mu N \longrightarrow N(x) = N_0 e^{-\mu x}$ For mono-energetic case, intensity is $I(\Delta x) = I_0 e^{-\mu \Delta x}$























X-Ray Imaging Equation At z = d there is no magnification, so $I_d(x,y) = I_0 \cos^3 \theta \cdot \exp\left(-\int_{L_{x,y}} \mu(s) ds / \cos \theta\right)$ $= I_0 \cos^3 \theta \cdot t_d(x,y)$ where $t_z(x,y)$ is the transmittivity of the object at distance z In general, with magnification $I_d(x,y) = I_0 \cos^3 \theta \cdot t_z(x/M(z), y/M(z))$

































Generalized Functions

Dirac delta functions are not ordinary functions that are defined by their value at each point. Instead, they are generalized functions that are defined by what they do underneath an integral.

The most important property of the Dirac delta is the sifting property $\int_{-\infty}^{\infty} \delta(x - x_0)g(x)dx = g(x_0) \text{ where } g(x) \text{ is a smooth function. This sifting}$ property can be understood by considering the limiting case $\lim_{n \to \infty} \int_{-\infty}^{\infty} \tau^{-1} \Pi(x/\tau)g(x)dx = g(x_0)$















